

Tools for Plasma Kinetics Simulation in PETSc

Matthew Knepley, Joseph Pusztaý, Mark Adams

Computer Science and Engineering
University at Buffalo

SIAM Computational Science and Engineering,
Cyberspace March 1, 2021



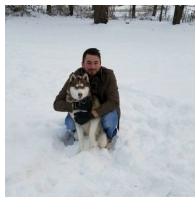
Collaborators

Plasma
Collaboration



Mark Adams

Projection and
Discrete Gradients



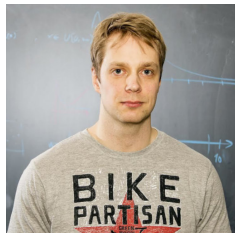
Joe Pusztay



Dan Finn

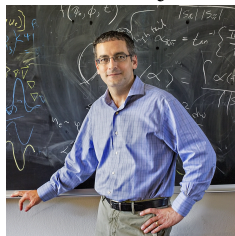
Collaborators

Metriplectic
Simulation



Eero Hirvijoki

Runaway
Electrons



Dylan Brennan

Outline

Structure-Preserving Evolution

Collisions

Vlasov Equation

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \frac{\partial f}{\partial \mathbf{x}} + \frac{d\mathbf{v}}{dt} \cdot \frac{\partial f}{\partial \mathbf{v}} = 0$$

We will need

- ▶ Particle discretization
- ▶ Symplectic integrators

Vlasov-Poisson Equation

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \frac{\partial f}{\partial \mathbf{x}} - \frac{q \nabla \phi}{m} \cdot \frac{\partial f}{\partial \mathbf{v}} = 0$$
$$-\Delta \phi = \rho$$

We will need

- ▶ FEM discretization
- ▶ FEM-Particle map

Weak Equivalence

$$f_{FEM} = \sum_i \phi_i \mathbf{f}_i \quad f_{PM} = \sum_p \delta(\mathbf{x} - \mathbf{x}_p) \delta(\mathbf{v} - \mathbf{v}_p) \mathbf{w}_p$$

Require that moments are preserved

$$\int \phi_i f_{FEM} = \int \phi_i f_{PM}$$
$$M \mathbf{f} = M_P \mathbf{w}$$

where M is the mass matrix, and

$$\begin{aligned} [M_P]_{ip} &= \int \phi_i \delta(\mathbf{x} - \mathbf{x}_p) \\ &= \phi_i(\mathbf{x}_p) \end{aligned}$$

Projections

Particle \rightarrow FEM (deposition)

$$Mf = M_P w$$

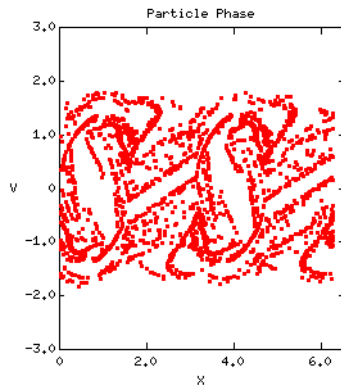
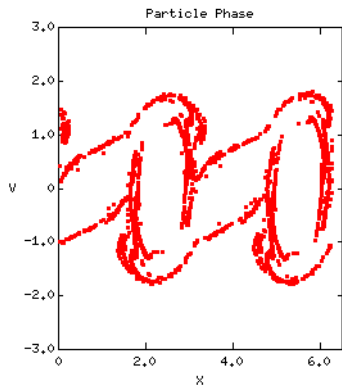
$$f = M^{-1} M_P w$$

FEM \rightarrow Particle

$$w = M_P^+ Mf$$

Examples

Two stream instability test in PETSc



(MollenAdamsKnepleyHagerChang2021)

Vlasov-Poisson-Landau Equation

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \frac{\partial f}{\partial \mathbf{x}} - \frac{q \nabla \phi}{m} \cdot \frac{\partial f}{\partial \mathbf{v}} = C(f)$$
$$-\Delta \phi = \rho$$

We will need

- ▶ Action of the Landau operator
- ▶ Metriplectic integrator

Long Time Evolution

Symmetries must be present in the discrete system:

- ▶ Conservation of mass
- ▶ Conservation of momentum
- ▶ Conservation of energy
- ▶ Monotonicity of entropy

Long Time Evolution

Tools we will need:

- ▶ Symplectic integrator will preserve moments
- ▶ Projection will preserve moments
- ▶ Discrete Landau must preserve moments
- ▶ DG integrator to preserve monotonic entropy
- ▶ Splitting should not destroy these properties

Discrete Gradients

Discrete Gradient integrators

$$\frac{f^{n+1} - f^n}{\Delta t} = S(f^n, f^{n+1}) \nabla F(f^n, f^{n+1})$$

satisfy a secant condition

$$(f^{n+1} - f^n) \cdot \nabla F(f^n, f^{n+1}) = F(f^{n+1}) - F(f^n).$$

Monotonicity of entropy means monotonicity of free energy:

$$\begin{aligned} S(f^{n+1}) - S(f^n) &= (E(f^{n+1}) - F(f^{n+1})) - (E(f^n) - F(f^n)) \\ &= - (F(f^{n+1}) - F(f^n)) \\ &= - (f^{n+1} - f^n) \cdot \nabla F \\ &= -\Delta t \nabla F \cdot S \cdot \nabla F \\ &\geq 0 \end{aligned}$$

since our S is symmetric negative semi-definite.

Outline

Structure-Preserving Evolution

Collisions

Collisions in the FE Basis

Collision in the Particle Basis

Outline

Collisions

Collisions in the FE Basis

Collision in the Particle Basis

Landau Operator

Strong Form

$$C_{\alpha}(f) = \sum_{\beta} \nu_{\alpha\beta} \frac{m_0}{m_{\alpha}} \nabla \cdot \int d\mathbf{v}' \mathbf{U}(\mathbf{v}, \mathbf{v}') \left(\frac{m_0}{m_{\alpha}} f_{\beta}(\mathbf{v}') \nabla f_{\alpha}(\mathbf{v}) - f_{\alpha}(\mathbf{v}) \nabla' f_{\beta}(\mathbf{v}') \frac{m_0}{m_{\beta}} \right)$$

$$\mathbf{U}(\mathbf{v}, \mathbf{v}') = \frac{1}{|\mathbf{v} - \mathbf{v}'|} \mathbf{I} - \frac{(\mathbf{v} - \mathbf{v}') \otimes (\mathbf{v} - \mathbf{v}')}{|\mathbf{v} - \mathbf{v}'|^3}$$

$$\nu_{\alpha\beta} = \frac{e_{\alpha}^2 e_{\beta}^2 \ln \Lambda_{\alpha\beta}}{8\pi m_0^2 \epsilon_0^2}$$

m_0 = reference mass

Landau Operator

Weak Form

$$\sum_{\beta} \nu_{\alpha\beta} \frac{m_0^2}{m_{\alpha}} \int d\mathbf{v}' \nabla' \psi(\mathbf{v}') \cdot \left(\frac{1}{m_{\alpha}} \mathbf{K}(f_{\beta}, \mathbf{v}') \phi(\mathbf{v}') + \frac{1}{m_{\beta}} \mathbf{D}(f_{\beta}, \mathbf{v}') \cdot \nabla' \phi(\mathbf{v}') \right)$$

$$\mathbf{K}(f, \mathbf{v}) = \int d\mathbf{v}' \mathbf{U}(\mathbf{v}, \mathbf{v}') \cdot \nabla' f(\mathbf{v}')$$

$$\mathbf{D}(f, \mathbf{v}) = \int d\mathbf{v}' \mathbf{U}(\mathbf{v}, \mathbf{v}') f(\mathbf{v}')$$

(Hirvijoki2017)

Landau Operator

Implementation

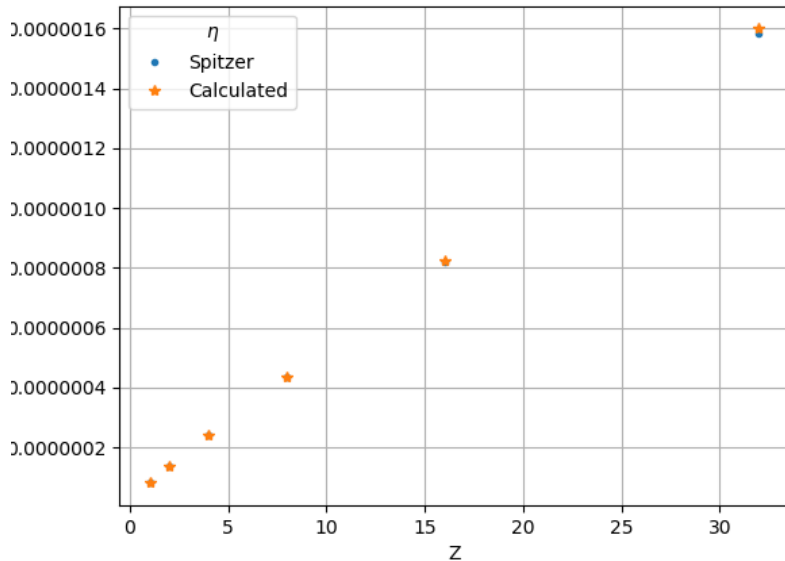
- ▶ Conservative projection to/from FEM
- ▶ Constant order Q_k
- ▶ Adaptive p4est grid

(AdamsHirvijokiKnepleyBrownIsaacMills2017)

Spitzer Resistivity

(AdamsKnepleyBrennan2021)

Spitzer vs calculated resistivity



Outline

Collisions

Collisions in the FE Basis

Collision in the Particle Basis

Particle Landau

(Hirvijoki2021)

If we ignore position, we have

$$\frac{d\mathbf{v}_p}{dt} = \frac{\nu}{m} \sum_{p'} w_{p'} \mathbf{Q}(\mathbf{v}_p - \mathbf{v}_{p'}) \Gamma(S, p, p')$$

where \mathbf{Q} is the Landau tensor,

$$\mathbf{Q}(\mathbf{v}) = \frac{1}{\|\mathbf{v}\|} \mathbf{I} - \frac{\mathbf{v} \otimes \mathbf{v}}{\|\mathbf{v}\|^3}$$

and

$$\Gamma(S, p, p') = \frac{1}{w_p} \frac{\partial S}{\partial \mathbf{v}_p} - \frac{1}{w_{p'}} \frac{\partial S}{\partial \mathbf{v}_{p'}}$$

$$S = - \int d\mathbf{v} f(\mathbf{v}) \ln f(\mathbf{v})$$

Particle Landau

(Hirvijoki2021)

If we ignore position, we have

$$\frac{d\mathbf{v}_p}{dt} = \frac{\nu}{m} \sum_{p'} w_{p'} \mathbf{Q}(\mathbf{v}_p - \mathbf{v}_{p'}) \Gamma(S, p, p')$$

where \mathbf{Q} is the Landau tensor,

$$\mathbf{Q}(\mathbf{v}) = \frac{1}{\|\mathbf{v}\|} \mathbf{I} - \frac{\mathbf{v} \otimes \mathbf{v}}{\|\mathbf{v}\|^3}$$

and

$$\Gamma(S, p, p') = \frac{1}{w_p} \frac{\partial S}{\partial \mathbf{v}_p} - \frac{1}{w_{p'}} \frac{\partial S}{\partial \mathbf{v}_{p'}}$$

$$S = - \sum_p w_p \ln w_p$$

Particle Landau

(Hirvijoki2021)

If we ignore position, we have

$$\frac{d\mathbf{v}_p}{dt} = \frac{\nu}{m} \sum_{p'} w_{p'} \mathbf{Q}(\mathbf{v}_p - \mathbf{v}_{p'}) \Gamma(S, p, p')$$

where \mathbf{Q} is the Landau tensor,

$$\mathbf{Q}(\mathbf{v}) = \frac{1}{\|\mathbf{v}\|} \mathbf{I} - \frac{\mathbf{v} \otimes \mathbf{v}}{\|\mathbf{v}\|^3}$$

and

$$\Gamma(S, p, p') = \frac{1}{w_p} \frac{\partial S}{\partial \mathbf{v}_p} - \frac{1}{w_{p'}} \frac{\partial S}{\partial \mathbf{v}_{p'}}$$

$$S_\epsilon = - \int d\mathbf{v} \sum_p w_p \psi_\epsilon(\mathbf{v} - \mathbf{v}_p) \ln \left(\sum_{p'} w_{p'} \psi_\epsilon(\mathbf{v} - \mathbf{v}_{p'}) \right)$$

Particle Landau

(Hirvijoki2021)

$$\frac{\mathbf{v}_p^{n+1} - \mathbf{v}_p^n}{\Delta t} = \frac{\nu}{m} \sum_{p'} w_{p'} \mathbf{Q} \left(\mathbf{v}_p^{n+1/2} - \mathbf{v}_{p'}^{n+1/2} \right) \Gamma(S_\epsilon^n, p, p')$$

where \mathbf{Q} is the Landau tensor,

$$\mathbf{Q}(\mathbf{v}) = \frac{1}{\|\mathbf{v}\|} \mathbf{I} - \frac{\mathbf{v} \otimes \mathbf{v}}{\|\mathbf{v}\|^3}$$

and

$$\Gamma(S, p, p') = \frac{1}{w_p} \frac{\partial S}{\partial \mathbf{v}_p} - \frac{1}{w_{p'}} \frac{\partial S}{\partial \mathbf{v}_{p'}}$$

$$S_\epsilon = - \int d\mathbf{v} \sum_p w_p \psi_\epsilon(\mathbf{v} - \mathbf{v}_p) \ln \left(\sum_{p'} w_{p'} \psi_\epsilon(\mathbf{v} - \mathbf{v}_{p'}) \right)$$

Future Work

- ▶ Prove that splitting preserves structure
- ▶ Scalability tests
- ▶ Higher order Discrete Gradients
- ▶ Mixed-Poisson FEM

References I