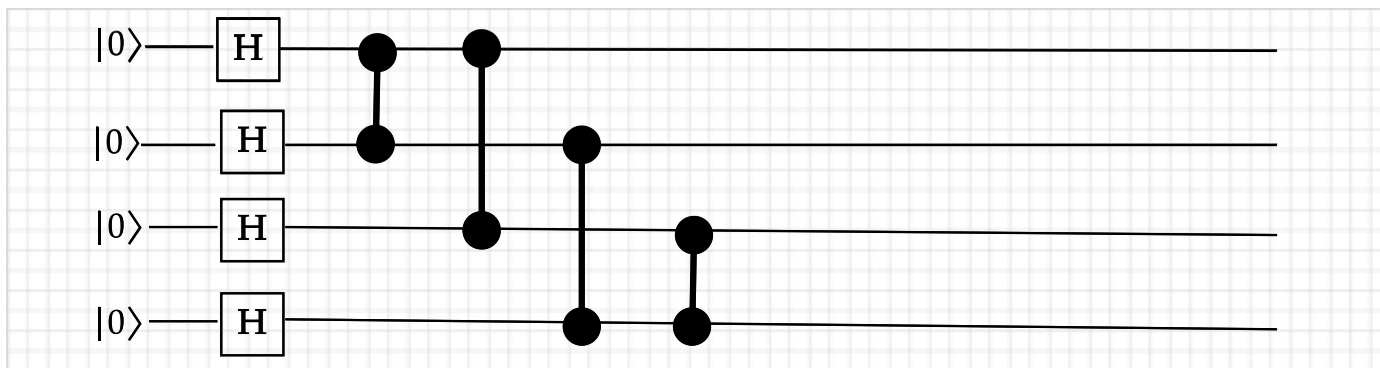


## CSE491/596 Lecture Mon. 11/20/23: General Quantum Circuits and Computations

If there are  $n$  qubits, then the underlying matrices we get are  $N \times N$  with  $N = 2^n$ . It is much harder to handle  $2^n$ -sized stuff than  $n$ -sized stuff. Happily, we can always break the basic gates down to constant size---3 at most with the Toffoli gate in practice---and there are theorems that guarantee constant size gates working in general. One important case of using  $n$  single-qubit gates is the **Hadamard transform**  $\mathbf{H} \otimes \mathbf{H} \otimes \dots \otimes \mathbf{H}$  ( $n$  times), which can be abbreviated  $\mathbf{H}^{\otimes n}$ :

$$H^{\otimes 2} = \frac{1}{2} \begin{bmatrix} 1 & 1 & 1 & 1 \\ 1 & -1 & 1 & -1 \\ 1 & 1 & -1 & -1 \\ 1 & -1 & -1 & 1 \end{bmatrix}, \quad H^{\otimes 3} = \frac{1}{2\sqrt{2}} \begin{bmatrix} 1 & 1 & 1 & 1 & 1 & 1 & 1 & 1 \\ 1 & -1 & 1 & -1 & 1 & -1 & 1 & -1 \\ 1 & 1 & -1 & -1 & 1 & 1 & -1 & -1 \\ 1 & -1 & -1 & 1 & 1 & -1 & -1 & 1 \\ \hline 1 & 1 & 1 & 1 & -1 & -1 & -1 & -1 \\ 1 & -1 & 1 & -1 & -1 & 1 & -1 & 1 \\ 1 & 1 & -1 & -1 & -1 & -1 & 1 & 1 \\ 1 & -1 & -1 & 1 & -1 & 1 & 1 & -1 \end{bmatrix}$$

We always have  $H^{\otimes n}|0^n\rangle = |+\rangle^{\otimes n} = |+\rangle^n =$  the all-1 vector of length  $N = 2^n$  divided by  $\sqrt{N} = \sqrt{2^n} = 2^{n/2}$ . Often this is the first step of a quantum circuit, for example:



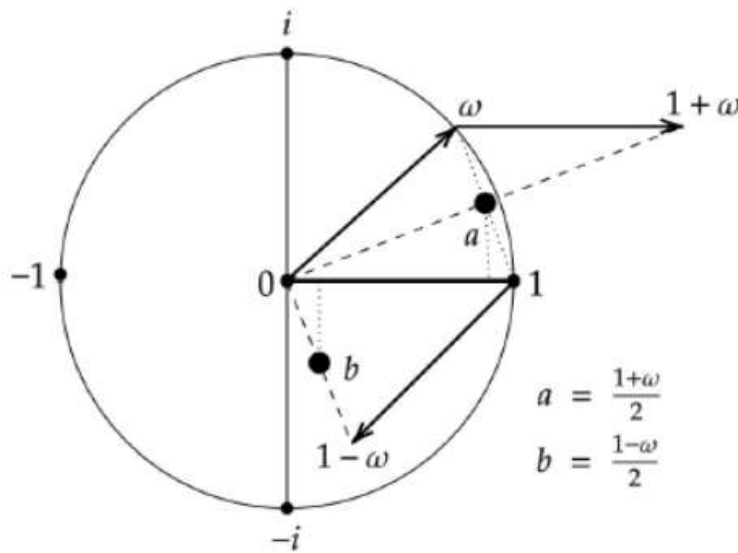
Putting the same Hadamard transform also at the end creates what is called a **graph state circuit**; we will analyze them later.

We will call an  $N \times N$  matrix that arises from a single small gate---or a tensor product of small gates---a **succinct** matrix. Thus a **quantum computation of length  $s$**  is formally a composition of  $s$  succinct matrices applied to some input vector. The text draws allusion to a classical computation on a binary string  $x$  of length  $n$ , such as  $x = 10100010$ , say. The quantum circuit starts with input the basis state  $|x\rangle = |10100010\rangle$ . We could actually start with  $|0^8\rangle$  but then **prepare** the state  $|x\rangle$  by making the first column of the circuit be the tensor product

$$\mathbf{X} \otimes \mathbf{I} \otimes \mathbf{X} \otimes \mathbf{I} \otimes \mathbf{I} \otimes \mathbf{I} \otimes \mathbf{X} \otimes \mathbf{I},$$

which has a NOT gate where  $x$  has a 1. This is why we often suppose ("without loss of generality") that the circuit starts with the all-zero basis vector.

The **Z** and **CZ** gates are the heads of an important family of basic gates having to do with rotations of **phase**, which is a curious but definitely physical property. When a complex number  $x + iy$  is rewritten in polar form as  $re^{i\theta}$ , the angle  $\theta$  is the phase. The magnitude is  $r$ , so when  $r = 1$  we have a unit complex number. Note that  $i$  itself is the same as  $e^{i\pi/2}$  since  $\frac{\pi}{2}$  means  $90^\circ$  phase. Then  $i^2 = e^{i\pi} = -1$  and if we put  $\omega = e^{i\pi/4}$  then  $\omega^2 = i$ . In Cartesian coordinates,  $\omega = \frac{1+i}{\sqrt{2}}$ . Here is some more geometry:



The vector  $\mathbf{u} = [a, b]^T$  is a funky unit vector. To see that it is a unit vector, note that

$$\|\mathbf{u}\|^2 = \langle \mathbf{u}, \mathbf{u} \rangle = \mathbf{u}^* \mathbf{u} = a^* a + b^* b = \left( \frac{1+\bar{\omega}}{2} \right) \left( \frac{1+\omega}{2} \right) + \left( \frac{1-\bar{\omega}}{2} \right) \left( \frac{1-\omega}{2} \right).$$

In polar form, the complex conjugate of  $e^{i\theta}$  is always  $e^{-i\theta} = e^{i(2\pi-\theta)}$ , so  $\bar{\omega} = e^{i7\pi/4} = \omega^7$ . In Cartesian coordinates,

$$\frac{1+\omega}{2} = \frac{1}{2} \left( 1 + \frac{1+i}{\sqrt{2}} \right) = \frac{\sqrt{2}+1+i}{2\sqrt{2}} \quad \text{and} \quad \frac{1-\omega}{2} = \frac{1}{2} \left( 1 - \frac{1+i}{\sqrt{2}} \right) = \frac{\sqrt{2}-1-i}{2\sqrt{2}}$$

So

$$\frac{1+\bar{\omega}}{2} = \frac{1}{2} \left( 1 + \frac{1-i}{\sqrt{2}} \right) = \frac{\sqrt{2}+1-i}{2\sqrt{2}} \quad \text{and} \quad \frac{1-\bar{\omega}}{2} = \frac{1}{2} \left( 1 - \frac{1-i}{\sqrt{2}} \right) = \frac{\sqrt{2}-1+i}{2\sqrt{2}}.$$

Then

$$\left(\frac{1+\bar{\omega}}{2}\right)\left(\frac{1+\omega}{2}\right) = \frac{1}{8}(\sqrt{2}+1+i)(\sqrt{2}+1-i) = \frac{1}{8}\left[(\sqrt{2}+1)^2 + 1\right] = \frac{1}{8}(2+1+2\sqrt{2}+1) = \frac{2+\sqrt{2}}{4}$$

and

$$\left(\frac{1-\bar{\omega}}{2}\right)\left(\frac{1-\omega}{2}\right) = \frac{1}{8}(\sqrt{2}-1-i)(\sqrt{2}-1+i) = \frac{1}{8}\left[(\sqrt{2}-1)^2 + 1\right] = \frac{1}{8}(2+1-2\sqrt{2}+1) = \frac{2-\sqrt{2}}{4}.$$

These squared values add to 1 as promised, so  $\mathbf{u} = [a, b]^T$  is a unit vector. How do we get it? Here is the start of an infinite family of gates:

$$\mathbf{Z} = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}, \quad \mathbf{S} = \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix}, \quad \mathbf{T} = \begin{bmatrix} 1 & 0 \\ 0 & \omega \end{bmatrix}, \quad \mathbf{T}_{\pi/8} = \begin{bmatrix} 1 & 0 \\ 0 & e^{i\pi/8} \end{bmatrix}.$$

The controlled versions to go with  $\mathbf{CZ}$  are  $\mathbf{CS}$ ,  $\mathbf{CT}$ , etc. They, too, are symmetric---indeed, all of these gates are controlled phase shifts conditioned on the basis-state 1 of all of the (one or two) qubits involved. (Here I must note global inconsistency and confusion in notation, especially about rotations, which we will try to resolve when we cover the **Bloch Sphere** next week.)

Now we have all the background we need to read **quantum circuits**. Lecture will go on to illustrate them, both out of section 4.5 and (the same examples) on QC web applets.