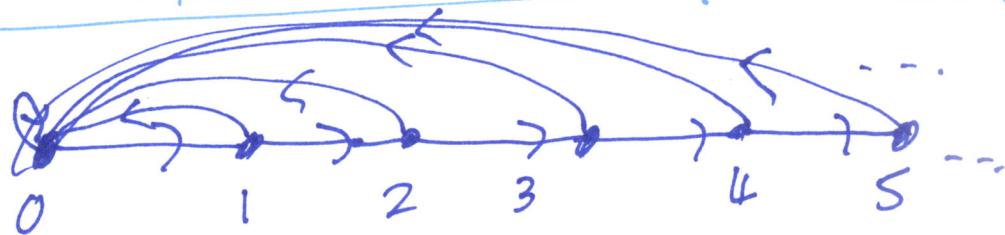


Why no good
randomized algm
is known for GAP
on directed graphs



Simplest example (IMHO) of a best-known randomized algm
involves not NL vs L, or polytime vs exp. time, but $\tilde{O}(n^2)$ vs $\tilde{O}(n)$

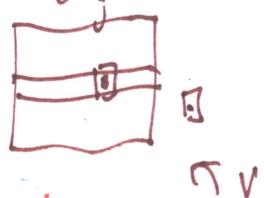
MATCHCHECK: INST: Three $n \times n$ matrices A, B, C (over some field F)
QUES: Does $A \cdot B = C$?

Deterministic: solve by multiplying $A \cdot B$, but this is simple way

Randomized: • Generate 1 or more vectors $v_1, \dots, v_r \in F^n$.

- For $j = 1 \dots r$, check whether $C v_j = A(B v_j)$.
- If any check fails, say no $\tilde{O}(n^2)$ time Repeat $\tilde{O}(n^2)$ time \therefore whole is $\tilde{O}(n^2)$ time
- If all checks succeed, then $\leq \frac{1}{2}r$ chance of getting a "bad v_j " each time

Abstractly, let $D = A \cdot B - C$. If $D \neq 0$ then D has at least one 1.



Then whether $v_j = 0$ or 1 flips the Parity of $Dv = A(Bv) - Cv$.

- We have a randomized algm R st. for all instances $x = \langle A, B, C \rangle$:
- $x \in \text{MATCHCHECK} \Rightarrow \Pr[\text{all checks succeed}] = 1 \Rightarrow \Pr[R \text{ is right}] = 1$.
- $x \notin \text{MATCHCHECK} \Rightarrow \Pr[\text{all succeed}] \leq \frac{1}{2}r \Rightarrow \Pr[R \text{ is wrong}] < \frac{1}{2}r$.

Defn: A language $L \subseteq \{0,1\}^*$ belongs to BQP [^{bounded error} _{Quantum poly time}] if there is a polynomial $p(n)$ and a quantum algorithm Q s.t. for all n and inputs $x \in \{0,1\}^n$:

$$x \in L \Rightarrow \Pr[Q(x) \text{ measures } 1 \text{ on qubit 1}] \geq \frac{3}{4} \quad \text{i.e. } \Pr[Q(x) \text{ is right}] \geq \frac{3}{4}$$

$$x \notin L \Rightarrow \Pr[Q(x) \text{ measures } 1 \text{ on like 1}] \leq \frac{1}{4} \quad \text{i.e. } \Pr[Q \text{ is wrong}] \leq \frac{1}{4}$$

i.e. $\Pr[Q(x) = L(x)]$ always $\geq \frac{3}{4}$. $\left\{ \begin{array}{l} \text{Can } \underline{\text{amplify}} \text{ "error" } \leq \frac{1}{4} \\ \text{to "error } \leq \frac{1}{r}" \text{ by } O(r) \text{ repeated trials.} \end{array} \right.$

The MATCHICK example, when time is poly vs exp not $n^2 \sqrt{n^3}$, typifies the definition of classical BPP. Thus $P \subseteq BPP \subseteq BQP$. Problems for which BQP may be larger include FACTORING, DISCRETE LOG, and $A \cdot B \stackrel{?}{=} C$ problems where A, B, C are exponentially big.

The standard model for universal quantum computing models Q as a family $[C_n]_{n=1}^\infty$ of quantum circuits using at least one gate per site. Clifford gates $H, I, X, Y, Z, S, \text{CNOT}, \text{CZ}$



00	01	10	11
01	10	00	00
01	01	10	00
10	00	10	00
11	00	01	-1

The three most common other gates are:



1	0
0	$\sqrt{2}$

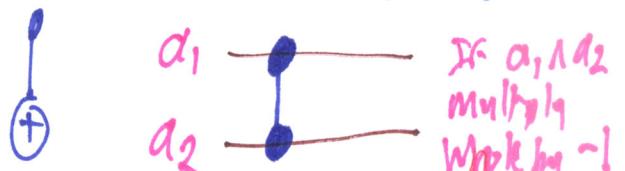
1	0	0	0
0	1	0	0
0	0	1	0
0	0	0	1

Toffoli

All non-Clifford.

T

CS



a_1

a_2

if $a_1 a_2$
multipl
w/ prob -1.

Recall $\mathbb{H} - \mathbb{B} - \mathbb{H}$? gives irrational probabilities for 1 and so BQP involves not simply coinflips.

The argument over BQP has two levels:

- ① Whether the Quantum Fourier Transform is really feasible to implement with $O(n^2)$ circuit of basic gates; see QFT_n Recursion below.

- ② Those gates involve "Higher T Twists": $\begin{bmatrix} 1 & 0 \\ 0 & \sqrt{i} \end{bmatrix}$

$$\begin{bmatrix} 1 & 0 \\ 0 & \sqrt{i} \end{bmatrix} \xrightarrow{\text{--- down to ---}} \boxed{TW_n} = \begin{bmatrix} 1 & 0 \\ 0 & \frac{1}{2}\sqrt{i} \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & i^{1/2^{n-2}} \end{bmatrix}$$

$TW_0 = I$
 $TW_1 = Z$
 $TW_2 = S$
 $TW_3 = T$

Do we only need to approximate such high precision?
exponentially small.

Added: Examples of the QFT. QFT_n takes $w = e^{\frac{2\pi i}{N}}$ where $N = 2^n$.

$n=1: N=2, w = e^{\frac{2\pi i}{2}} = e^{\pi i} = -1$. $\boxed{QFT_n[i,j] = w^{i,j} \quad i,j \in 0, \dots, N-1}$

Row 0 and Column 0 are always all 1s. So $QFT_1 = \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} = H_1$ Normalize $\frac{1}{\sqrt{2}}$.

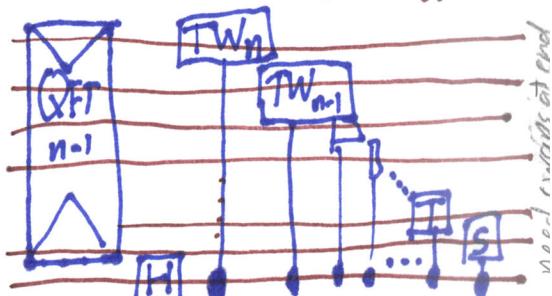
$n=2, N=4, w = e^{\frac{2\pi i}{4}} = e^{\pi i/2} = i$,
Then $w^2 = -1, w^3 = -i, w^4 = 1$.

$n=3, N=8, w = e^{\frac{2\pi i}{8}} = \sqrt[8]{i}$

1	1	1	1	1	1	1	1
1	w	i	w ³	-1	w ⁵	-i	w
1	i	-1	-i	1	i	-1	-i
1	w ³	-i	w	-1	w ⁷	i	w ⁵
1	-1	1	-1	1	-1	1	-1
1	w ⁵	i	w	-1	w ⁷	-i	w ³
1	-i	-1	i	1	-i	-1	i
1	w	-i	w ⁵	-1	w ³	i	w

$QFT_2 = \begin{bmatrix} 1 & 1 & 1 & 1 \\ 1 & i & -1 & -i \\ 1 & -1 & 1 & -1 \\ 1 & -1 & -1 & i \end{bmatrix} \cdot \frac{1}{2}$

As with the Hadamard transform, the normalizing constant is $\frac{1}{\sqrt{N}} = 2^{-n/2}$.



Important Fact: The state $x = (0.5, \frac{i}{2}, -0.5, \frac{-i}{2})$ equals $a_0|000\rangle + a_1|01\rangle + a_2|10\rangle + a_3|11\rangle$, where $a_0 = \frac{1}{2}, a_1 = \frac{i}{2}, a_2 = -\frac{1}{2}, a_3 = -\frac{i}{2}$ as a vector over \mathbb{C}^4 i.e. $a_0e_0 + a_1e_1 + a_2e_2 + a_3e_3$ the standard basis. All quantum circuits C, being linear transformations, obey the linear rule

$$Cx = [C](a_0e_0 + a_1e_1 + a_2e_2 + a_3e_3) = a_0(Ce_0) + a_1(Ce_1) + a_2(Ce_2) + a_3(Ce_3) = \sum_{j=0}^3 a_j(C|j\rangle)$$

The behavior of C on any input state is a linear function of its behavior on basis states, see binarising